

MPO 663 - Convective and Mesoscale Met.
Homework 2

The point of this problem set is to illustrate the *Rossby radius of deformation*, and its dependence on vertical structure. We use the simplest possible framework for a heating-induced balanced vortex: a hydrostatic Boussinesq stratified atmosphere, with perturbations linearized about a state of rest (no advection terms). Thus, $\rho = \text{constant}$ (except that there is a buoyancy b), so that the pressure gradient force is written as $\phi_x \equiv p'_x / \rho$. (I used π instead of ϕ in class.)

Please **understand** our equation set:

$u_t = fv - \phi_x$	u momentum equation
$v_t = -fu - \phi_y$	v momentum equation
$0 = b - \phi_z$	w momentum equation (hydrostatic)
$u_x + v_y + w_z = 0$	continuity equation
$b_t = -N^2 w + q$	thermodynamic equation w/ heating ($q = g/\theta \, d\theta/dt$)

We study solutions with separable vertical and horizontal structure. For this step, it is simplest to go to 2 dimensions (ignore y variations) and to set aside the Coriolis term, which can be trivially added back in later. This simplifies the equation set to 4.

Understand them:

$u_t = -\phi_x$	u momentum equation
$0 = b - \phi_z$	w momentum equation (hydrostatic)
$u_x + w_z = 0$	continuity equation
$b_t = -N^2 w + q$	thermodynamic equation w/ heating q

1. Assume solutions of the form $u = U(x,t) \cos(mz)$, $\phi = \Phi(x,t) \cos(mz)$, $w = W(x,t) \sin(mz)$, $b = B(x,t) \sin(mz)$, $q = Q(x,t) \sin(mz)$.

- a) (i) **Confirm** that these satisfy the boundary condition $w=0$ at $z=0$ for all vertical wavenumbers m .

At $z=0$, $w=W(x,t)\sin(m(0))=W(x,t)[0]=0$

(ii) **Plug** these forms into the 4 equations above,

$$u_t = -\phi_x \leftrightarrow U_t(x, t) = -\Phi_x(x, t)$$

$$0 = b - \phi_z \leftrightarrow 0 = B(x, t) + m\Phi(x, t)$$

$$u_x + w_z = 0 \leftrightarrow U_x(x, t) + mW(x, t) = 0$$

$$b_t = -N^2 w + q \leftrightarrow B_t(x, t) = -N^2 W(x, t) + Q(x, t)$$

(iii) then **substitute** the hydrostatic equation into the thermo equation

$$-m\Phi_t(x, t) = -N^2 W(x, t) + Q(x, t)$$

(iii) and the result into the continuity equation to **show**:

$$U_t = -\Phi_x \quad \text{u momentum equation}$$

$$\Phi_t = -N^2/m^2 (U_x + \delta_d) \quad \text{height equation w/ source}$$

where the heating is expressed as *adiabatic divergence* $\delta_d \equiv \partial/\partial z(q/N^2) = -mQ/N^2$.

$$U_t(x, t) = -\Phi_x(x, t) \quad \text{from (i) above}$$

$$W(x, t) = \frac{1}{N^2} (m\Phi_t(x, t) + Q(x, t))$$

$$U_x + mW = U_x(x, t) + \frac{m}{N^2} (m\Phi_t(x, t) + Q(x, t)) = 0$$

$$\Phi_t(x, t) = \frac{N^2}{m^2} \left(-U_x(x, t) + \frac{m}{N^2} Q(x, t) \right) = -\frac{N^2}{m^2} \left(U_x(x, t) - \frac{mQ(x, t)}{N} \right)$$

$$\Phi_t(x, t) = -\frac{N^2}{m^2} (U_x(x, t) + \delta_d(x, t))$$

b) (i) Rewrite these to make clear that they are the equations for forced waves with free-surface height amplitude $h = \Phi/g$, in a shallow (compared to wavelength) layer of water (or any other inviscid fluid), with *equivalent depth* $H = N^2/(gm^2)$.

Plugging in $\Phi=hg$ and $gH = N^2/m^2$ into the two highlighted equation above gives:

$$U_t = -(gh)_x = -gh_x \quad \text{(u momentum equation)}$$

$$\Phi_t = -gH(U_x + \delta_d) \quad \text{(height equation w/ source)}$$

(ii) **Confirm** that the v equation, Coriolis terms, and u and v momentum source terms could have been carried along in the same way to give the forced shallow water set for a given value of m:

$$\Phi_t = -gH(U_x + \delta_d)$$

$$U_t = fV - \Phi_x + F \quad \text{u momentum equation}$$

$$V_t = -fU - \Phi_y + G \quad \text{v momentum equation}$$

$$\Phi_t = -N^2/m^2 (U_x + V_y + \delta_d) \quad \text{how divergence changes } \Phi$$

(i)

First, Let $u = U(x,y,t) \cos(mz)$, $v = V(x,y,t) \cos(mz)$,
 $\phi = \Phi(x,y,t) \cos(mz)$, $w = W(x,y,t) \sin(mz)$,
 $b = B(x,y,t) \sin(mz)$, $q = Q(x,y,t) \sin(mz)$

At $z=0$, $w=W(x,y,t)\sin(m(0))=W(x,y,,t)[0]=0$

(ii)

$$u_t = fv - \phi_x + X \leftrightarrow U_t(x,y,t) = fV(x,y,t) - \Phi_x(x,y,t) + F$$

$$v_t = -fu - \phi_y + Y \leftrightarrow V_t(x,y,t) = -fU(x,y,t) - \Phi_y(x,y,t) + G$$

$$0 = b - \phi_z \leftrightarrow 0 = B(x,y,t) + m\Phi(x,y,t)$$

$$u_x + v_y + w_z = 0 \leftrightarrow U_x(x,y,t) + V_y(x,y,t) + mW(x,y,t) = 0$$

$$b_t = -N^2w + q \leftrightarrow B_t(x,y,t) = -N^2W(x,y,t) + Q(x,y,t)$$

Note: X, Y, F, and G represent friction.

(iii)

$$-m\Phi_t(x,y,t) = -N^2W(x,y,t) + Q(x,y,t)$$

(iiii)

$$W(x,y,t) = \frac{1}{N^2} (Q(x,y,t) + m\Phi_t(x,y,t))$$

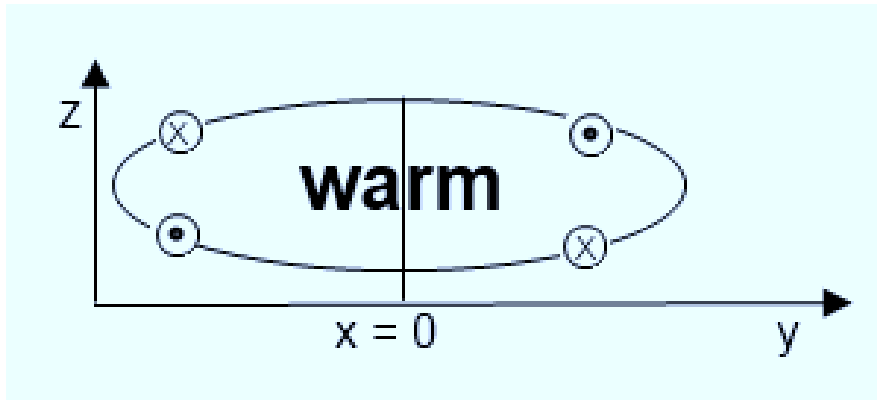
$$U_x + V_y + mW = U_x(x,y,t) + V_y(x,y,t) + \frac{m}{N^2} (m\Phi_t(x,y,t) + Q(x,y,t))$$

$$= 0; \quad \Phi_t(x,y,t) = \frac{N^2}{m^2} \left(-U_x(x,y,t) - V_y(x,y,t) + \frac{m}{N^2} Q(y,x,t) \right)$$

$$= -\frac{N^2}{m^2} \left(U_x(x,y,t) + V_y(x,y,t) - \frac{mQ(x,y,t)}{N} \right)$$

$$\Phi_t(x,y,t) = -\frac{N^2}{m^2} (U_x(x,y,t) + V_y(x,y,t) + \delta_d(x,y,t))$$

2. Now let's solve for the structure of an inviscid ($F=G=0$) balanced line vortex (shear line) created by a *past* heating along $x = 0$ on an f-plane, with $f = 10^{-4} \text{ s}^{-1}$. Suppose that the heating which created the vortex was a line of heating along $x=0$, and only excited a single vertical wavenumber m . Sketch this flow in the x - z plane for $m=1/2$ and z domain $[0,1]$. (It is like the cross section of a hurricane: a warm core vortex).



a) **Confirm** that this balanced flow should have no y or t derivatives (i.e., the flow is steady state) with $\delta_d = 0$ at the current time.

-Since the heat source at $x = 0$ was infinitely long in the y -direction, there is no y -derivative in the current flow.

-Since the heating occurred at some time in the past, all the gravity waves are already gone and the flow is steady-state in time, and thus there is no t -derivative. Also, since the heating occurred in the past, $\delta_d = -mQ/N^2 = 0$ at the current time.

b) (i) **Write** the steady state version of the 3-equation set above.

$$U_t(x, y, t) = fV(x, y, t) - \Phi_x(x, y, t) + F \rightarrow \Phi_x(x) = fv(x)$$

$$V_t(x, y, t) = -fU(x, y, t) - \Phi_x(x, y, t) + G \rightarrow \Phi_x(x) = -fu(x)$$

$$\Phi_t(x, y, t) = -\frac{N^2}{m^2} (U_x(x, y, t) + V_y(x, y, t) + \delta_d(x, y, t)) \rightarrow u_x(x) = 0$$

(ii) **What, in words, remains of the equation set's physical meaning?**

Geostrophic balance

(iii) Does it give us the structure of the vortex?

No.

(iii) Does it contain the information that we specifically seek the vortex *created by a past heating along $x = 0$* ? Why not?

No. While the geostrophic balance obtained above is valid, this result is ambiguous, as the content/information of the thermodynamic equation with heating is lost when the y and t derivatives are cancelled, plus the height equation,

$$\Phi_t = -gH(U_x + V_y + \delta_d), \text{ becomes just } 0 = 0.$$

3. To proceed, we need a different tactic. **Derive** the potential vorticity conservation equation for the 2D shallow water set from 1b (with $F=G=Q=0$), by **constructing** a vorticity equation from U and V and **substituting** divergence from the height equation.

$$u_t = fv - \phi_x, v_t = -fu - \phi_y, u = U(x,y,t) \cos(mz), v = V(x,y,t) \cos(mz),$$

$$\phi = \Phi(x,y,t) \cos(mz), \zeta = v_x - u_y$$

$$U_t = fV - \Phi_x, V_t = -fU - \Phi_y$$

$$\zeta_t = (V_x - U_y)_t = (V_t)_x - (U_t)_y = (-fU_x - \Phi_{yx}) - (fV_y - \Phi_{xy})$$

$$\zeta_t = -f(U_x + V_y) \quad \text{vorticity equation}$$

$$h = \Phi/g; \quad \Phi = gh; \quad \Phi_t = -gH(U_x + V_y + \delta_d); \quad \delta_d = 0$$

$$gh_t = -gH(U_x + V_y); \quad h_t = -H(U_x + V_y) \quad \text{height equation}$$

$$\frac{\zeta_t}{f} = -(U_x + V_y); \quad \frac{\zeta_t}{f} = \frac{h_t}{H} \quad \text{PV equation}$$

4.

a) To define the balanced vortex structure, **express mathematically** the statement that PV is constant and equal to its pre-heating value everywhere except along the line $x=0$ where the heating occurred.

$$\frac{\zeta_t}{f} = \frac{h_t}{H} = \frac{f}{H}; \quad \zeta_t = \frac{f^2}{H} = \frac{fh_t}{H} = \text{constant in } t \rightarrow h_t = \text{constant in } t$$

$$\rightarrow \frac{\zeta}{f} = \frac{h_t}{H}(t) = \frac{h}{H} \quad (\text{PV equation, time integrated})$$

$$V = V_g = \frac{1}{f} \frac{\partial \Phi}{\partial x} = \frac{g}{f} h_x \quad (V \text{ is geostrophic})$$

$$V_x = \frac{g}{f} h_{xx}; \quad \frac{\zeta}{f} = \frac{V_x - U_y}{f}; \quad U_y = 0; \quad \frac{\zeta}{f} = \frac{V_x}{f} = \frac{h}{H}; \quad V_x = \frac{g}{f} h_{xx} = \frac{hf}{H}$$

$$h_{xx} = \frac{f^2}{gH} h$$

$$\text{Guess: } h = h_0 e^{-\frac{f}{\sqrt{gH}}x} \quad (\text{picked the decaying solution, so that } h_x < 0)$$

$$\text{Proof: } h_x = -h_0 \frac{f}{\sqrt{gH}} e^{-\frac{f}{\sqrt{gH}}x}; \quad h_{xx} = \frac{f^2}{gH} \left(h_0 e^{-\frac{f}{\sqrt{gH}}x} \right) = \frac{f^2}{gH} h \quad \text{Confirmed!}$$

So, $PV = \frac{\zeta}{f} = \frac{h}{H}$ is constant and equal to its preheating value of $\frac{h_0}{H}$ along $x = 0$.

b) **Solve** this equation for $\Phi(x)$ and $V(x)$. **Write** the horizontal length scale for the vortex's exponential decay structure: this length is called the *Rossby deformation radius*. What is its dependence on m ?

$$L_R = \frac{\sqrt{gH}}{f}; \text{ Let } c = \sqrt{gH} = \text{phase speed, then } L_R = c/f$$

$$H = \frac{N^2}{gm^2}; \quad c = \sqrt{g \left(\frac{N^2}{gm^2} \right)} = \frac{N}{m}$$